

**What information can you get from this technique?**



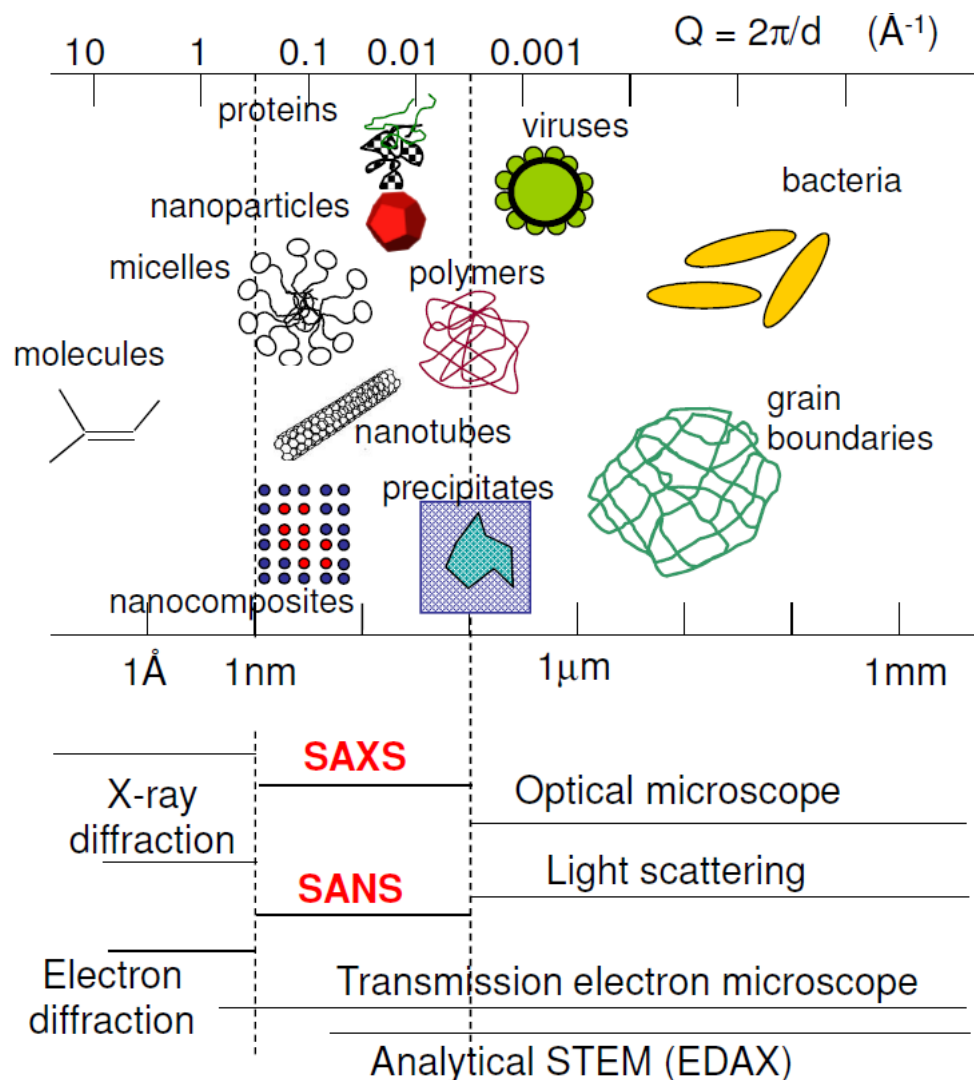
## A wide range of fields:

- Medicine
- Biology
- Chemistry
- Physics
- Archaeology
- Environmental and conservation sciences
- Materials

## A wide range of systems:

- Polymer processing
- Self assembly of mesoscopic metal particles
- Colloids
- Inorganic aggregates
- Liquid crystals
- The supramolecular organisation in biological systems
- The structure and function of muscle filaments
- Corneal transparency

# Size range comparisons



Wavelength (Energy) range	0.9 - 1.9 Å ( 6.5 -13 keV)
Flux at sample	>1.25 10 <sup>11</sup> ph/s 1.24 Å for a beam current of 250 mA
Bandpass ( $\Delta E/E$ )	< 10 <sup>-4</sup>
Beam size at sample	Variable between ~65 - 1200 $\mu\text{m}$ horizontally ~30 - 265 $\mu\text{m}$ vertically
Beam divergence at sample	<0.5x0.1 mrad <sup>2</sup>
Q range SAXS	0.0066-0.7 Å <sup>-1</sup>
2 $\theta$ range WAXS	3.0° - 62°

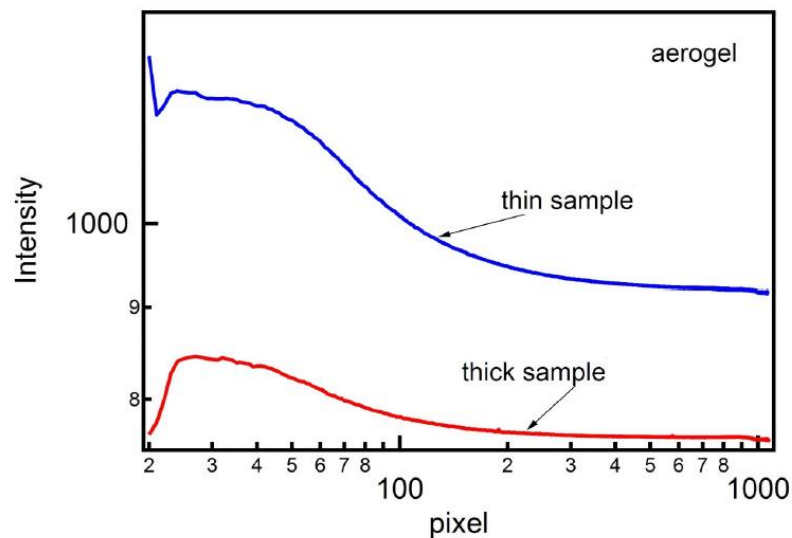
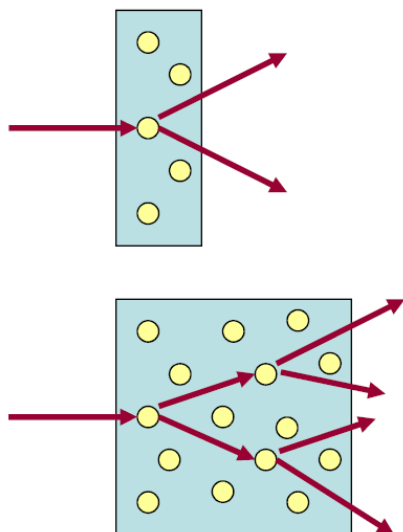
- Thermo-stated liquid cell (Required volume on demand)
- Non thermo-stated liquid cell rack for 24 samples (Volume 30ml)
- Thermo-stated ladder for 6 capillaries (1mm or 2mm diameter)
- Non thermo-stated ladder for 20 capillaries (1mm diameter)
- One film holder for 45 samples
- 1 Linkam stage for 1mm capillaries
- 1 Linkam stage for 22mm diameter films

	Controller	Pumping system	Linkam Stage	
Sample type			Capillary	Film
Model	T95	LNP954	HFSX350-CAP	THMS600
Max Temp (°C)	1500	-196	350	600
Max rate (°C/min)	200	100	30	150

## Thickness of the sample

- Affect transmission
- Affect the shape of the curve (Difficult to analyse)

Aim to 70% transmission



## Thickness of the sample

The maximum scattering intensity is achieved by selecting the optimal thickness of the sample.

The scattering intensity can be formulated as:

$d$ : the thickness of the sample

$\mu$ : linear absorption coefficient

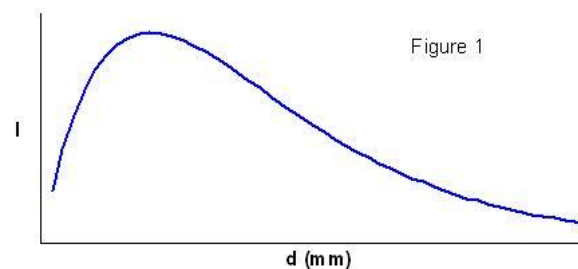


Figure 1

Optimal thickness 
$$\frac{\partial I}{\partial d} = k(1 - \mu d)e^{-\mu d} \Rightarrow d_{optimal} = \frac{1}{\mu}$$

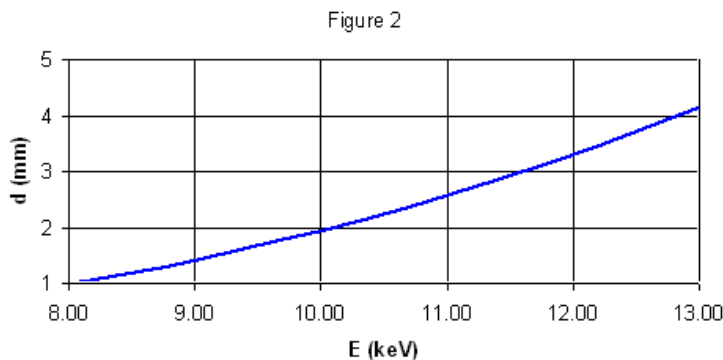


Figure 2

The linear absorption coefficient can be calculated or measured.

Optimal thickness of the sample (water in this case) as a function of the energy.

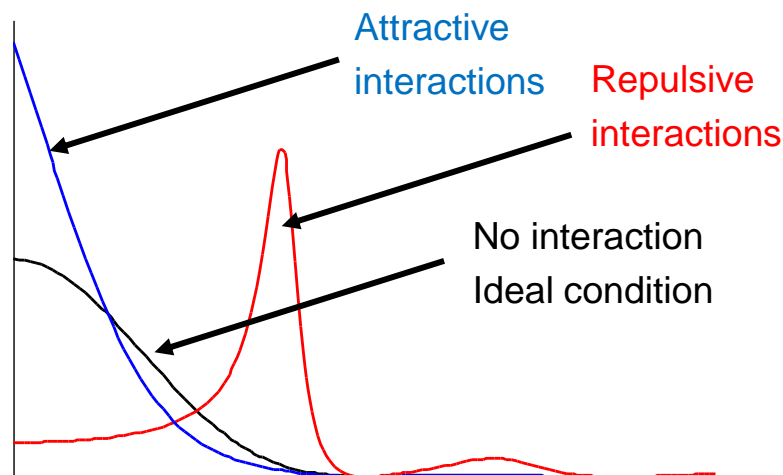
## Concentration

Big particle scatter more

Higher concentration

better signal **But don't burn out the detector**

can complicate analysis



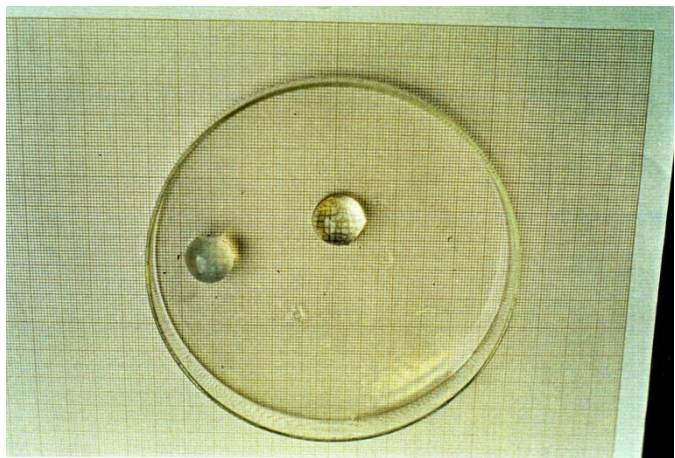
Minimum concentration for synchrotron:  $\approx 1\text{mg/ml}$



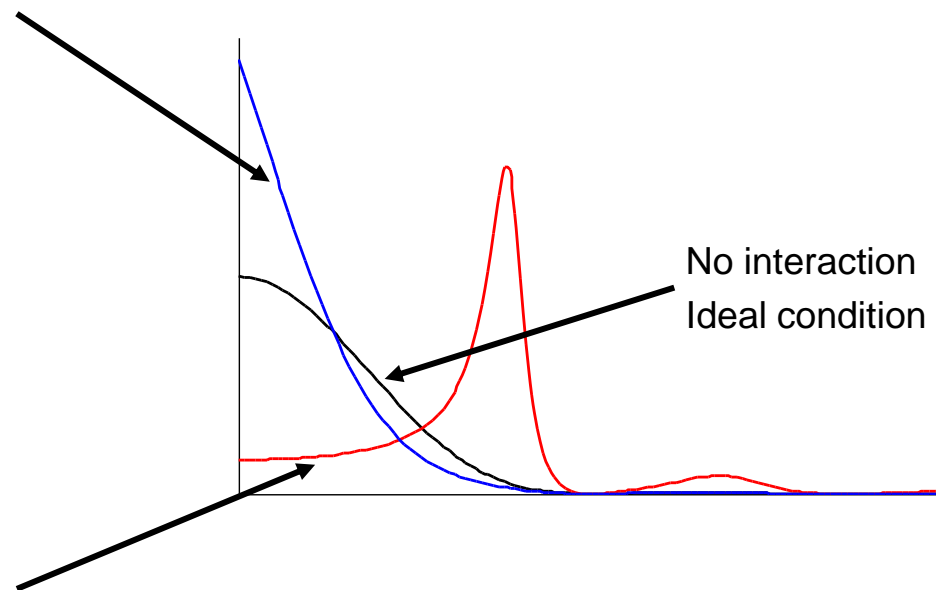
## High concentration Interaction between particles



Attractive  
interactions



Repulsive  
interactions

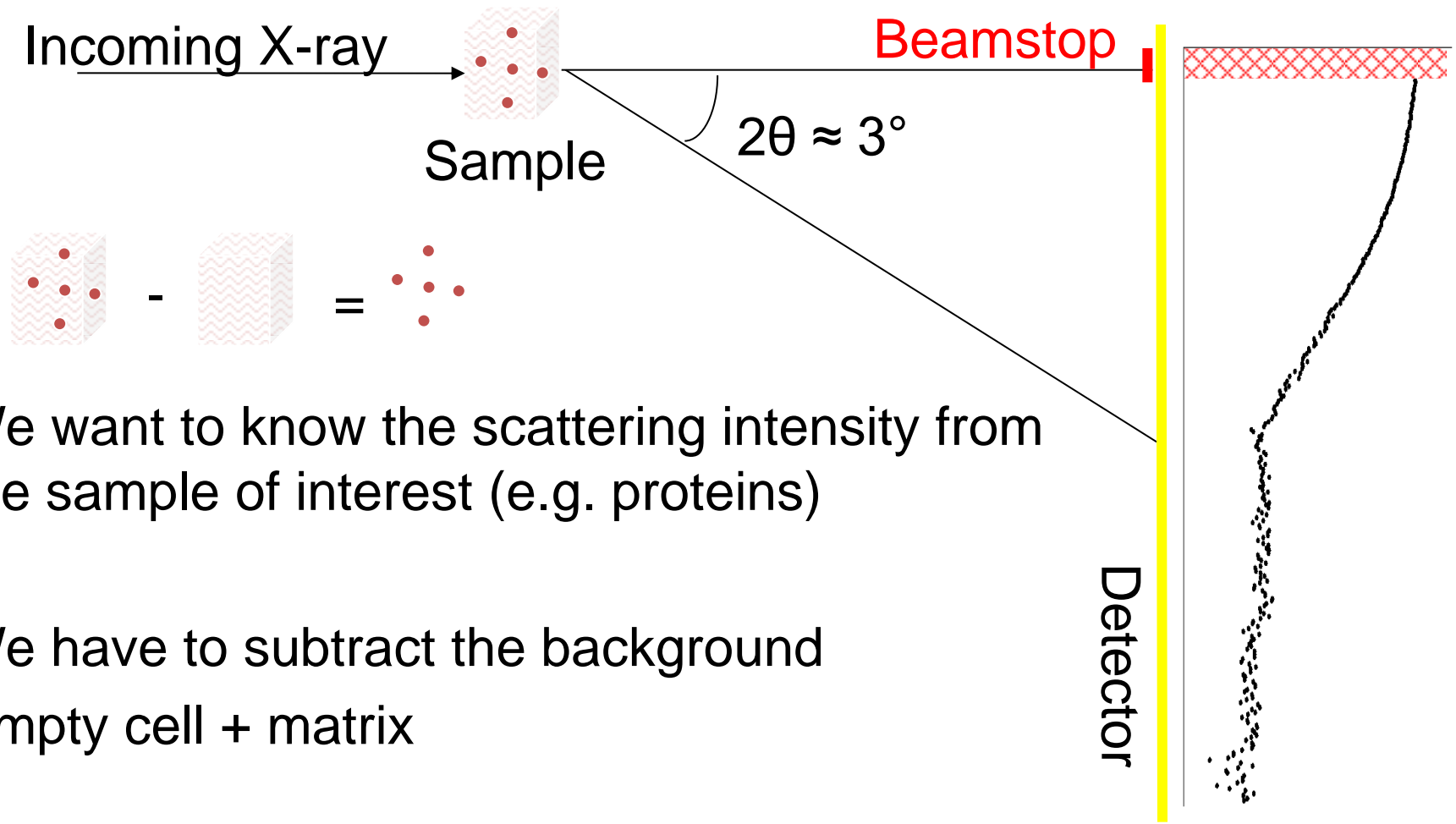


## Contrast



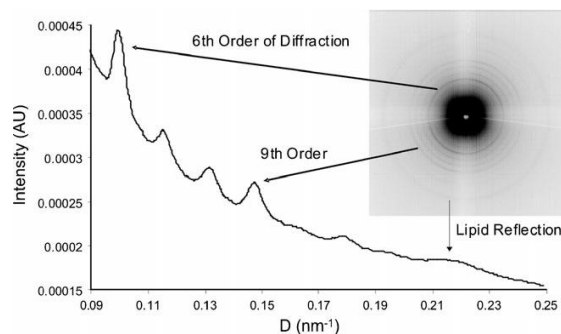
When the monster came, Lola, like the peppered moth remained motionless and undetected. Harold, of course, was immediately devoured.

The electronic density of the particle **MUST** not match the electronic density of the matrix.



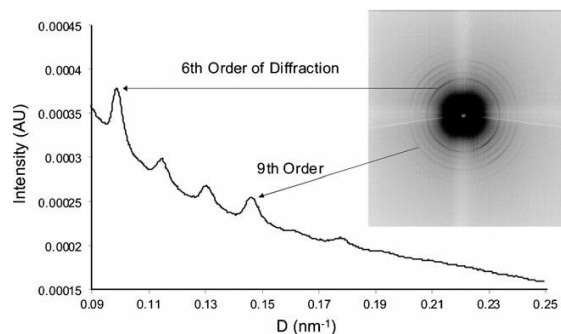
We want to know the scattering intensity from the sample of interest (e.g. proteins)

We have to subtract the background  
Empty cell + matrix



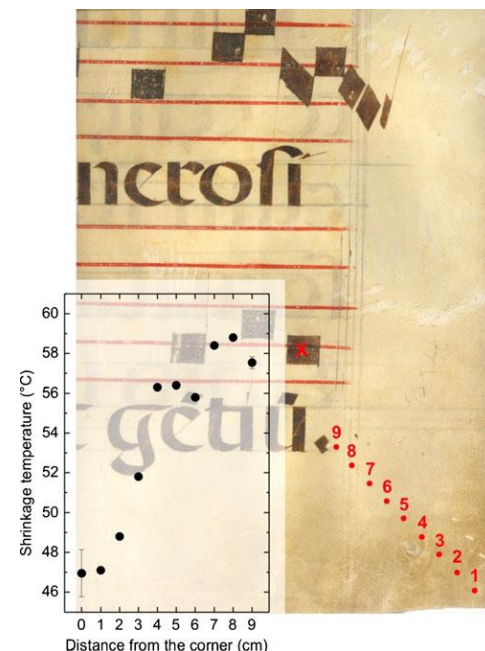
After 1 h of extraction, the  $0.22 \text{ nm}^{-1}$  reflection is still observed. However, after 4 h of extraction, the  $0.22 \text{ nm}^{-1}$  reflection no longer appears, indicating that the lipid phase of the parchment sample has been removed.

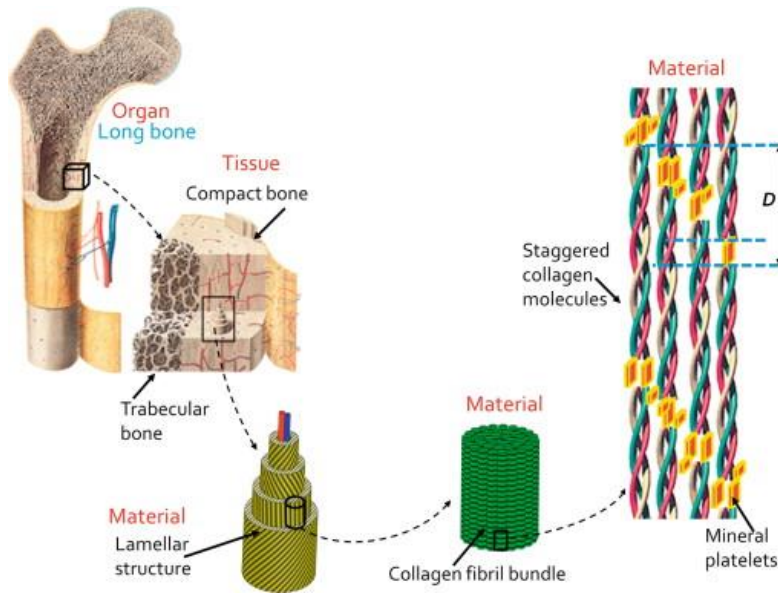
Lipids decrease shrinkage temperature.



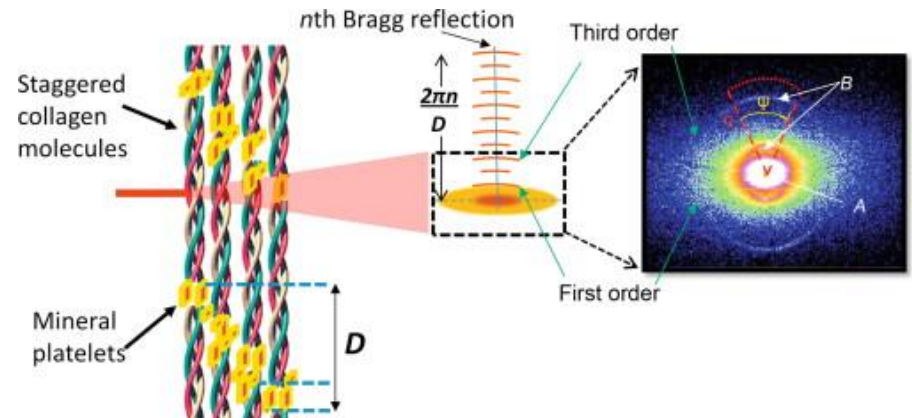
Parchment more degraded at the corner of the page by manual handling.

Možir *et al.*, A study of degradation of historic parchment using small-angle X-ray scattering, synchrotron-IR and multivariate data analysis, *Anal Bioanal Chem* (2012) 402:1559–1566 DOI 10.1007/s00216-011-5392-6





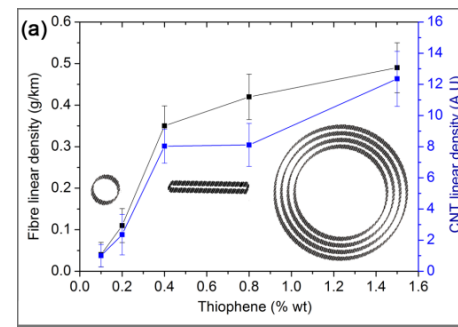
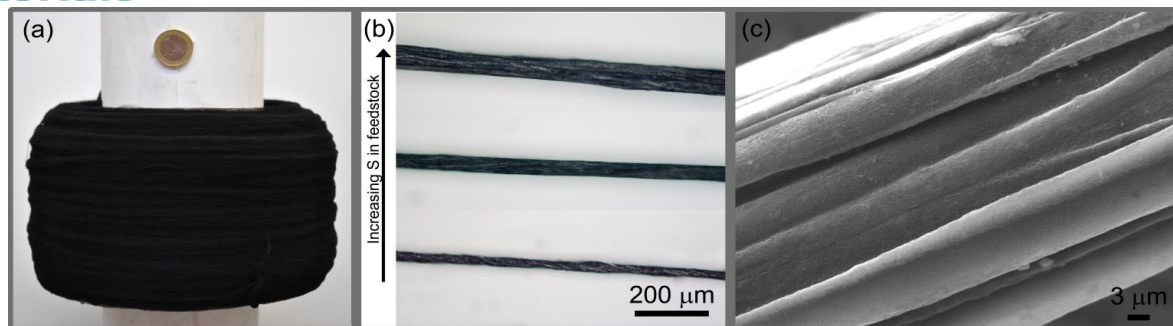
Strain, stress, and other mechanical parameters determination at small scales ( $< 100$  nm) in nanostructured biomineralized composites.





# Carbon nanotube - 1-

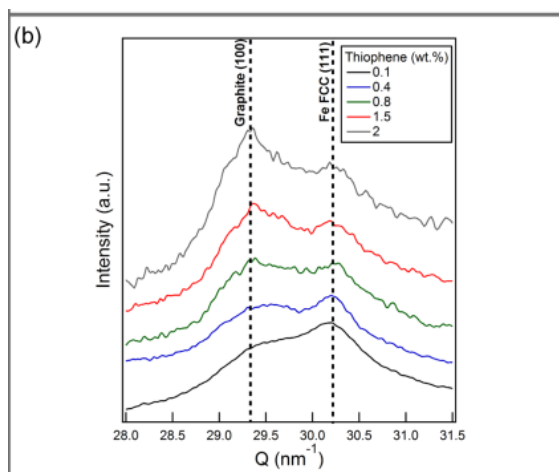
The research group has a process to produce kilometres of continuous macroscopic fibres of CNTs



**Synchrotron XRD** confirms the increase of graphitic layers at turbostratic separation

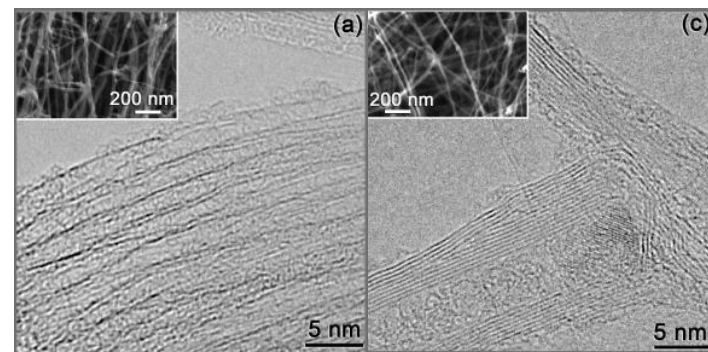
$S \uparrow$  in feedstock  $\rightarrow$  Intensity ratio  $[(100)_{\text{Graphite}} / (111)_{\text{Fe}}]$

They can tailor the type of nanotubes through the addition of sulphur precursor (thiophene)



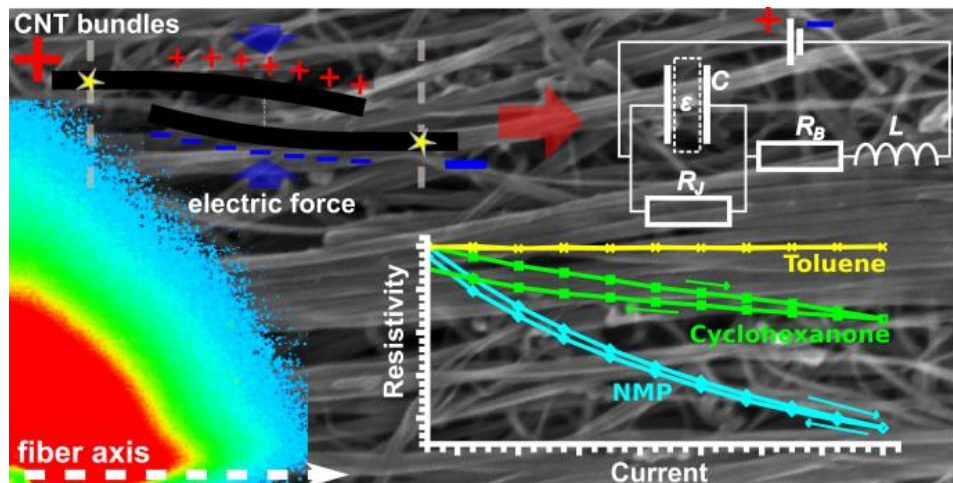
Low S  $\rightarrow$  SWCNTs

High S  $\rightarrow$  MWCNTs

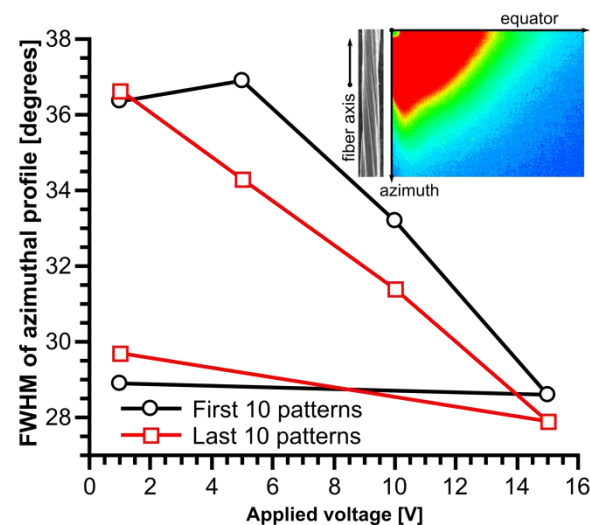


Reguero *et al.*, Controlling carbon nanotube type in macroscopic fibers synthesized by the direct spinning process, 2014, Chemistry of Materials, **26** (11), 3550-3557 DOI: 10.1021/cm501187x

The CNT fibres have a complex hierarchical structure. Its mesoporosity makes liquids infiltrate the fibre, which changes its electrical conductivity.

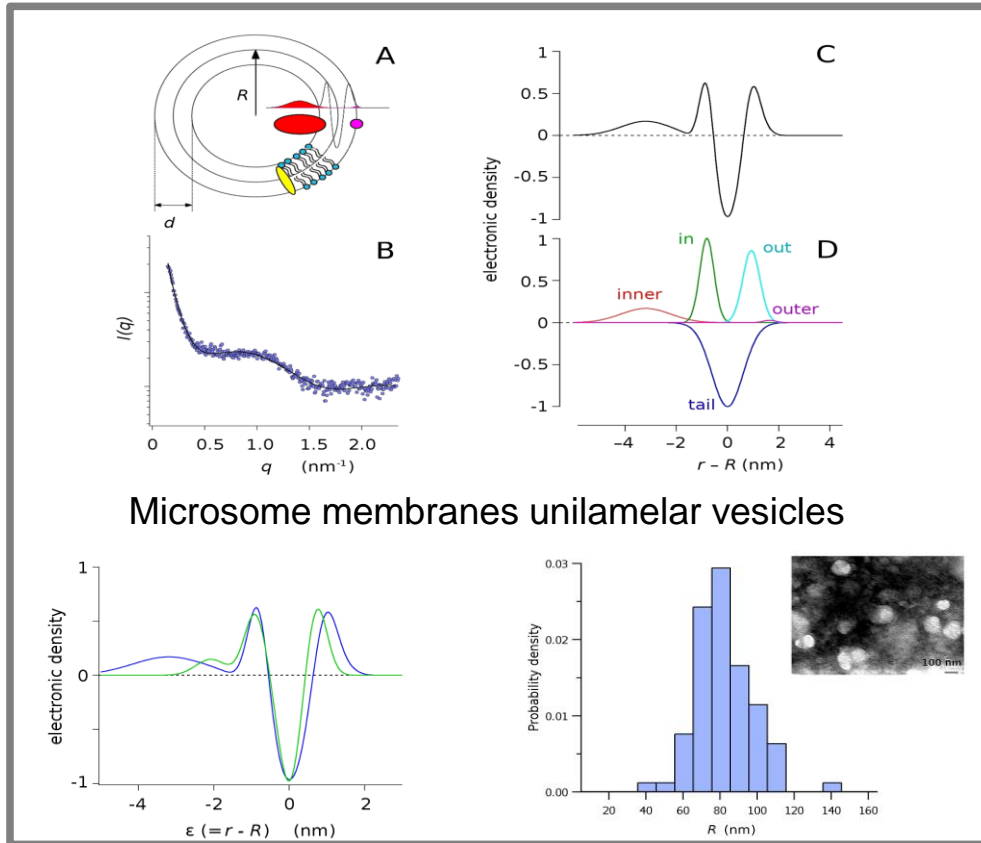


This non-Ohmic behaviour is a manifestation of nanoscale effects observed on a macroscopic scale. It arises due to fibre swelling, which can be observed by SAXS.



Terrones *et al.*, Electric Field-Modulated Non-ohmic Behavior of Carbon Nanotube Fibers in Polar Liquids, 2014, ACS NANO, **8** (8) , 8497-8504, DOI: 10.1021/nn5030835

From scattering theory, analytic expressions are derived for the bilayer form factor over a spherical geometry, assuming the lipid bilayer electron density to be composed of a series of Gaussian shells.



A) Sketch of the model of a vesicle consistent with the measured SAXS data.

B) SAXS spectra of the WT-CFTR membranes

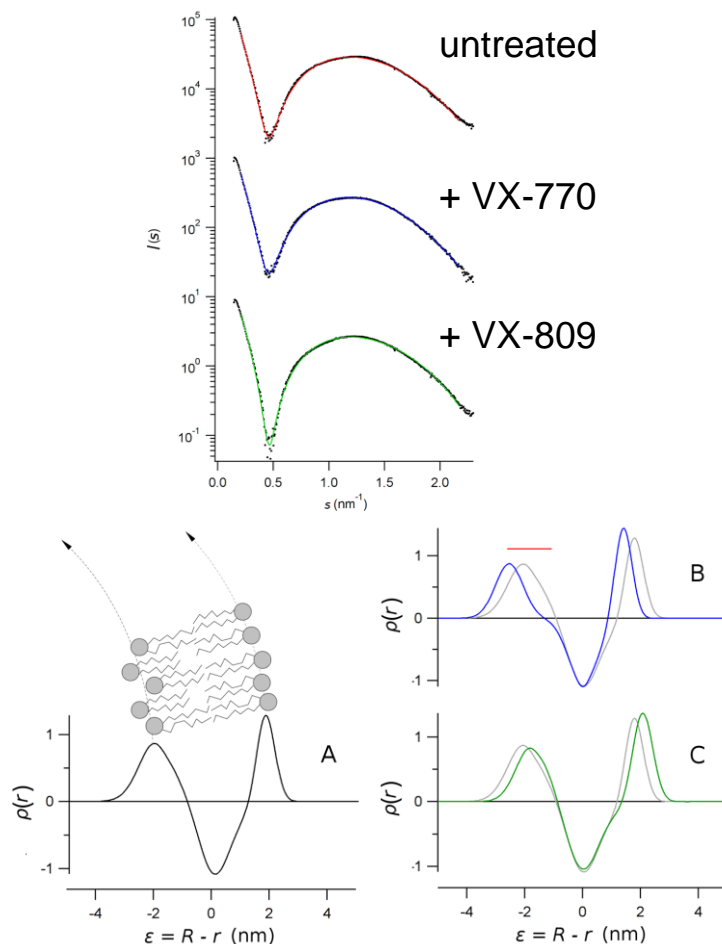
C) The electronic density profile was calculated from the five Gaussian model

D) The decomposition of each singular Gaussian used to model the electronic density

Electron density profiles of from microsome vesicles wall of native cells (green) and cells overexpressing CFTR (blue; protein involved in cystic fibrosis). Micrographs of vesicles are shown.

Baroni et al., Direct interaction of a CFTR potentiator and a CFTR corrector with phospholipid bilayers, European Biophysics Journal, July 2014, Volume 43, Issue 6-7, pp 341-346





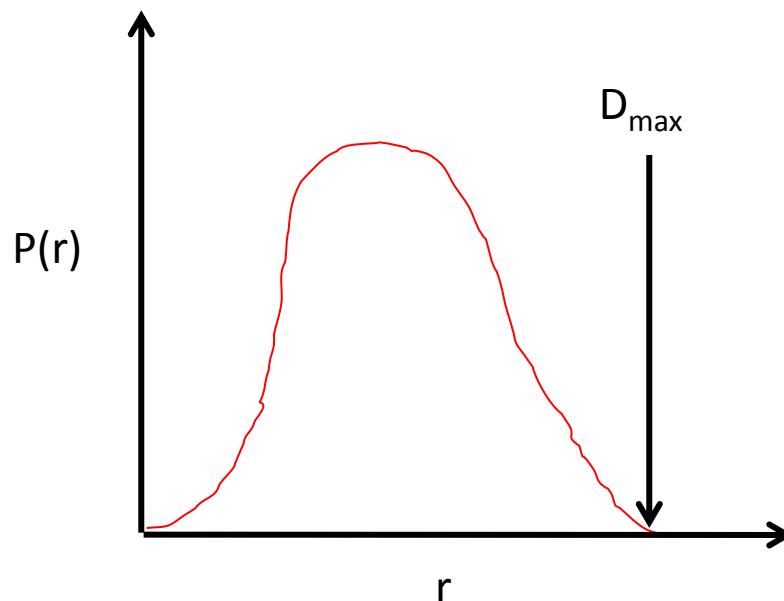
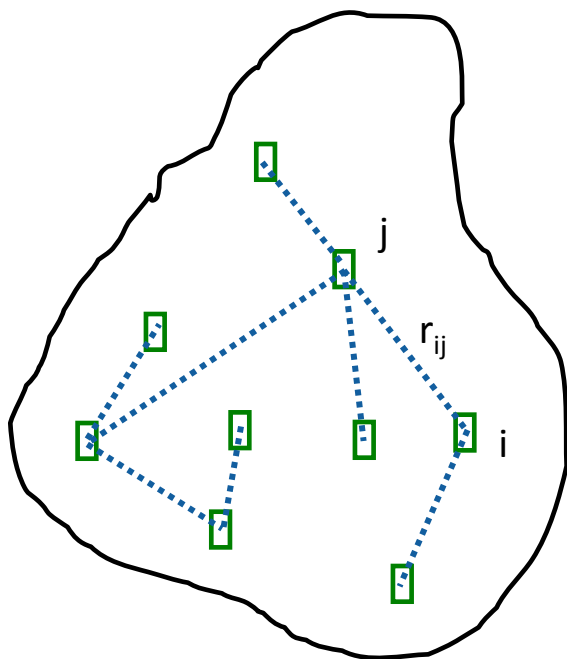
Large unilamellar liposomes (LUV) obtained by extrusion of phospholipids, treated with drugs for the cystic fibrosis therapy.

Electron density of the LUVs wall showing the destabilization induced by drugs.

$$I(q) = \int_{V_r} \int_{V_{r'}} \Delta\rho(r) \Delta\rho(r') e^{-i\vec{q}(\vec{r}-\vec{r}')} dV_r dV_{r'}$$

$$p(r) = \frac{1}{(2\pi)^2} \int_0^\infty I(q) q r \sin(qr) dq$$

Probability of finding a point at  $r$  from a given point.



$$p(r) = \frac{1}{(2\pi)^2} \int_0^\infty I(q) q r \sin(qr) dq$$

In theory, very easy calculation

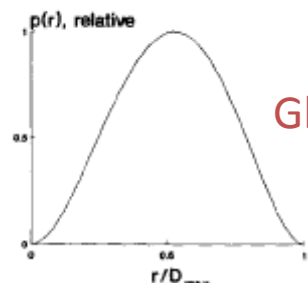
Problem:

The intensity :

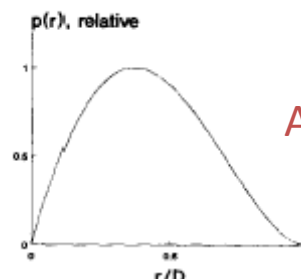
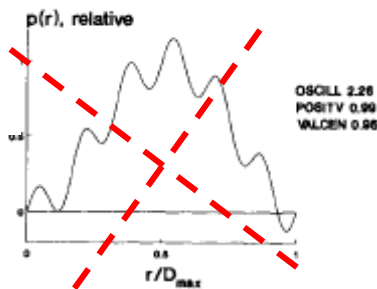
- ✓ only known over  $q_{\min}$ - $q_{\max}$  (Detector size)
- ✓ affected by experimental errors

⇒ Fourier transform of incomplete and noisy data is a ill-posed problem

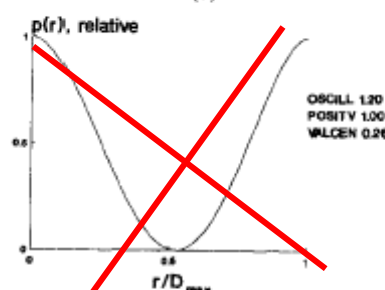
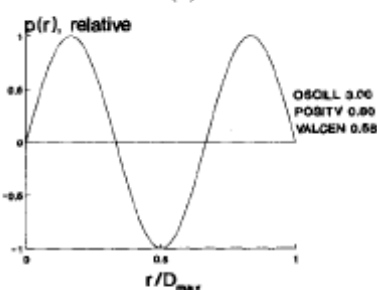
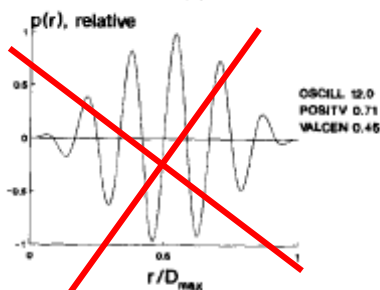
**Solution: Indirect Fourier Transform**



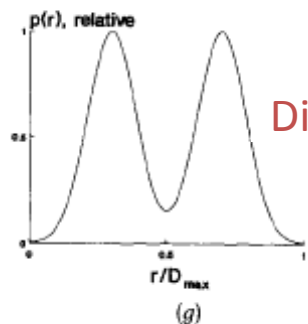
Globular



Anisotropic particle



Bilayer



Dimers

*J. Appl. Cryst.* (1992). **25**, 495–503

## Determination of the Regularization Parameter in Indirect-Transform Methods Using Perceptual Criteria

By D. I. SVERGUN\*†

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# Particle in solution - 4 -

## Guinier Law

$$I(q) = 4\pi \int_0^\infty p(r) \frac{\sin(qr)}{qr} dr$$

Taylor series expansion of

$$\frac{\sin(qr)}{qr} = 1 - \frac{(qr)^2}{3!} + \frac{(qr)^4}{5!} - \dots$$

$$I(q) = \underbrace{4\pi \int_0^\infty p(r) dr}_{I(0)} \left[ 1 - \underbrace{\left( \frac{q^2}{3!} \frac{\int_0^\infty r^2 p(r) dr}{\int_0^\infty p(r) dr} \right)}_{2R_g^2} \right] = I(0) \left[ 1 - \frac{q^2 R_g^2}{3} \right]$$

Taylor series expansion of  $e^x = 1 + x + \frac{x^2}{2!} + \frac{x^3}{3!} + \dots$  with  $x = -\frac{q^2 R_g^2}{3}$

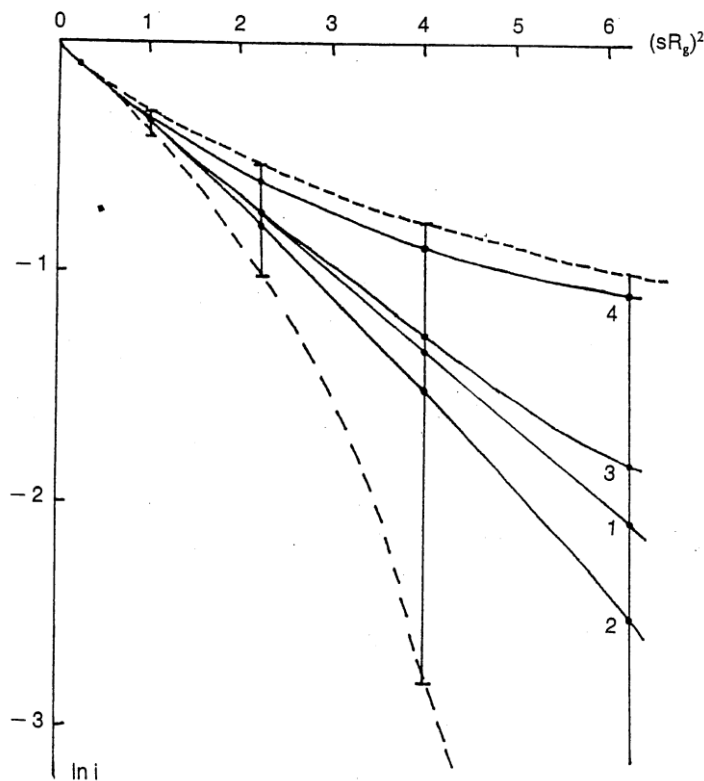
$$qR_g \ll 1$$

$$I(q) = I(0) e^{-\frac{q^2 R_g^2}{3}}$$

# Particle in solution - 5 -

## Accuracy of the Guinier Law

Valide for  $qR_g \ll 1$



1 – Guinier law (exponential)

2 – Sphere

3 – Thin disk

4 – long rod

Depends on the shape of the particle

$qR_g=1.5$  about 20.0% – 30.0% error

$qR_g < 1.3$  to reduce errors

*Structure Analysis by Small Angle X-ray and Neutron Scattering*  
L.A. Feigin and D.I. Svergun (1987), Plenum Press.

## Radius of gyration

$$R_g^2 = \frac{\int_{V_r} \Delta\rho(r) r^2 dV_r}{\int_{V_r} \Delta\rho(r) dV_r}$$

$R_g$  is the quadratic mean of distances to the centre of mass weighted by the contrast of electron density

$R_g$  is an **index of non sphericity**

$$\text{Sphere: } R_g = \sqrt{\frac{3}{5}} R \quad \text{Cylinder (D,H) } R_g = \sqrt{\frac{D^2}{8} + \frac{H^2}{12}}$$

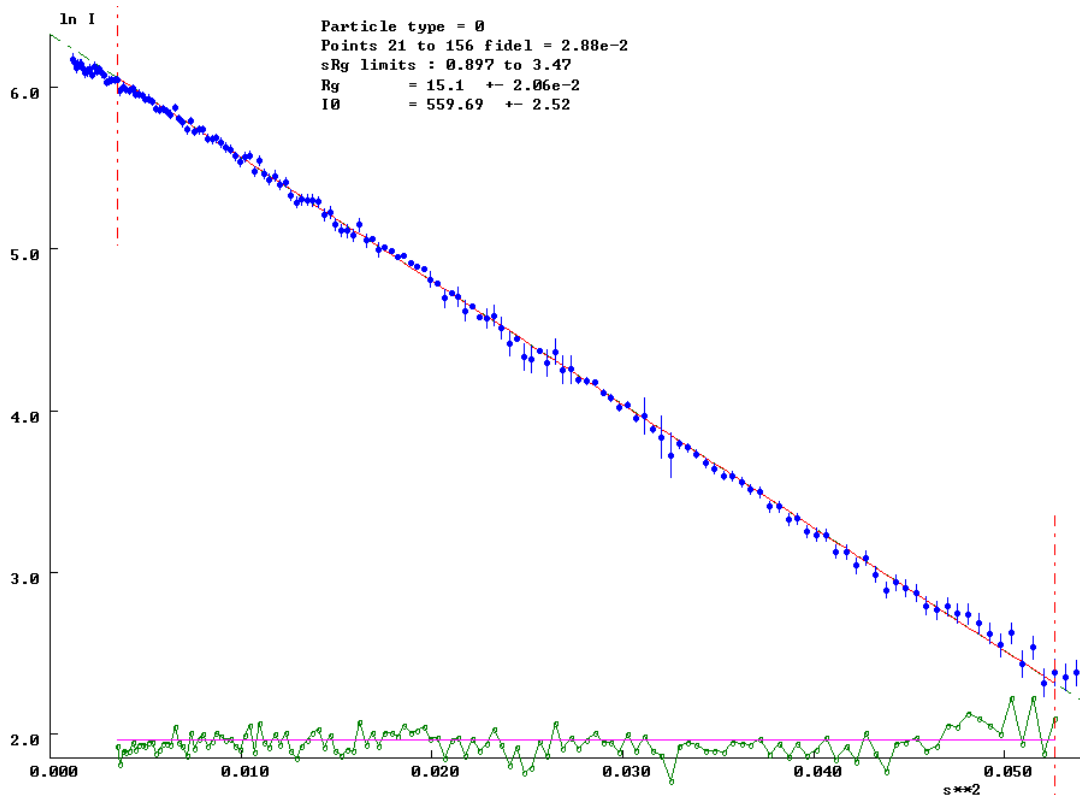
Smallest  $R_g$  for a given volume

$$\text{Ellipsoid of revolution (a,b) } R_g = \sqrt{\frac{2a^2 + b^2}{5}}$$

# Particle in solution - 7 - Guinier Plot

$$I(q) = I(0)e^{-\frac{q^2 R_g^2}{3}}$$

$$\ln I(q) = \ln I(0) - \frac{q^2 R_g^2}{3}$$

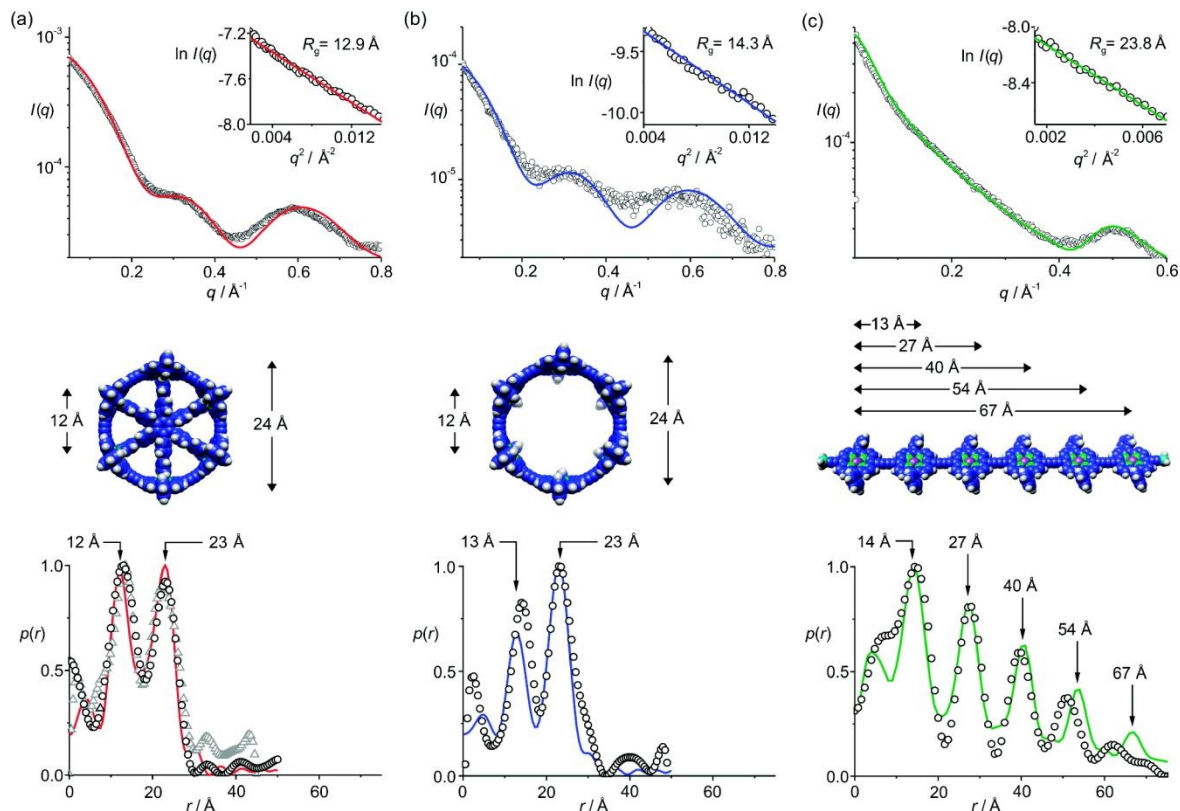


Linear regression

- ✓ Slope ->  $R_g^2$
- ✓ Intercept ->  $I(0)$   
->  $M_w$



Peak position = Zinc position



O'Sullivan *et al.*, Vernier templating and synthesis of a 12-porphyrin nano-ring, *Nature*, (2011) 469, 72–75

Beamline Responsible:

Marc Malfois

Beamline scientist:

Christina Kamma-Lorger

Juan-Carlos Martinez

Postdoctoral Research Associate:

Eva Crosas

Technician:

Jordi Prat

Controls:

Gabriel Jover

Engineering:

Joaquin Gonzalez

Carles Colldelram

Electronics:

Abel Fontseré

