

Structural insights into the barocaloric effect in $\text{Fe}^{\text{II}}(\text{pap-5NO}_2)_2$ spin-crossover complex.

Dra. Gloria Subías Peruga

David Gracia, Vera Cuartero, Marco Evangelisti, Javier Blasco, Emmanouil Charkiolakis, Catalin Popescu, Marie-Laure Boillot and Talal Mallah

1. Motivation

- Barocaloric green refrigeration
- Why molecular Fe^{II} spin-crossover complexes?

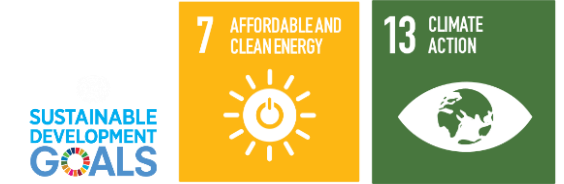
2. Results

- i. Low-pressure synchrotron powder X-ray diffraction
- ii. Calorimetry

3. Conclusions

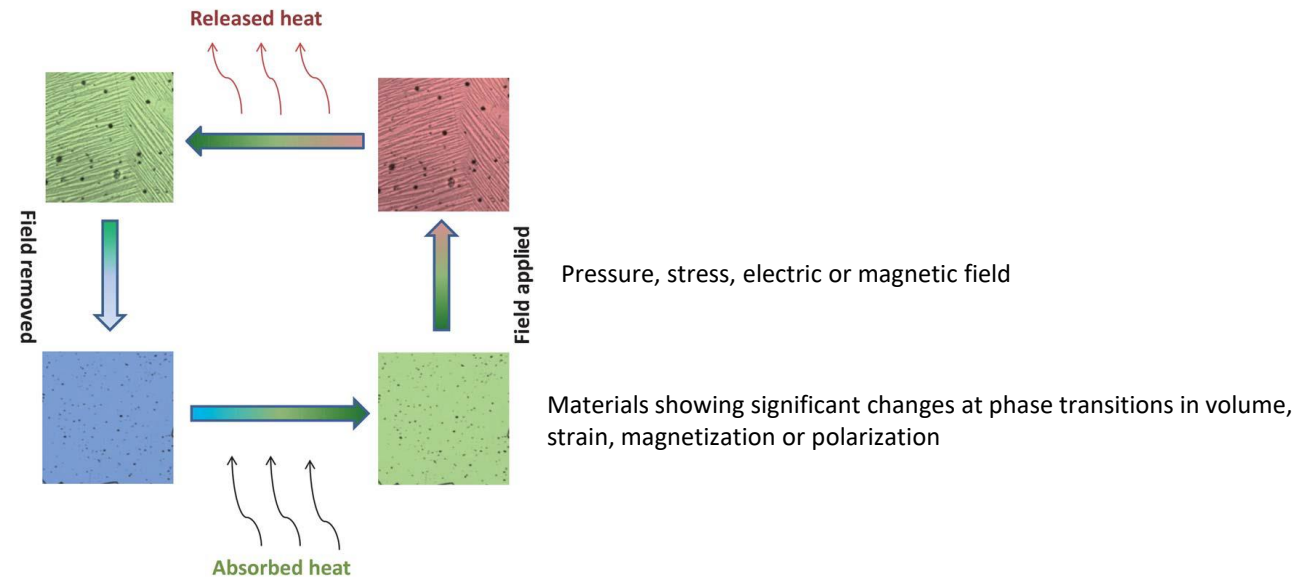
Motivation – Barocaloric green refrigeration

- **Refrigeration** is **essential** for *health and comfort, conservation of food and medicines, electronics ...* but it represents 30% of total electricity and it is a growing energy sector.
- **Vapor-compression**, the current technology:
 - ☞ HFC-based refrigerant gases → global warming
 - ☞ Limited energy efficiency (40%)
- **Solid-state caloric refrigeration**, very promising alternative:
 - ☞ Climate-friendly solid materials
 - ☞ Larger energy efficiency (60%)



Magnetocaloric (MC), elastocaloric (eC),
Electrocaloric (EC), **barocaloric (BC)** effects

Barocaloric effects in the solid state,
edited by P. Lloveras. IOP Publishing 2023



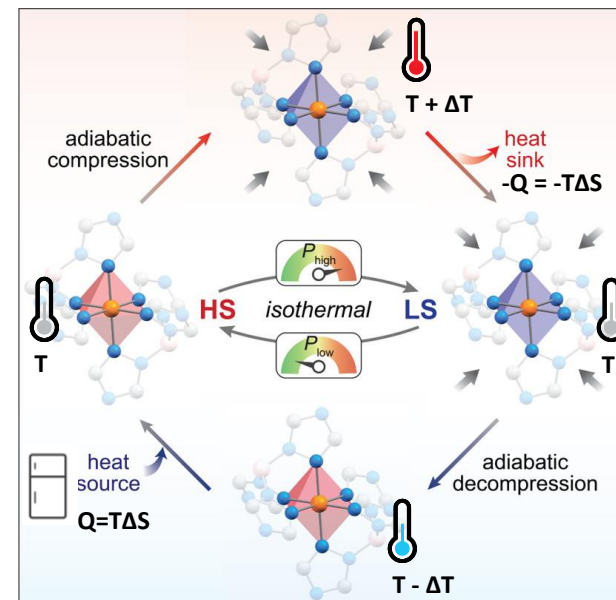
L. Mañosa et al., J. Mater. Chem. A, 2013,1, 4925-4936

Motivation – Why molecular Fe^{II} spin-crossover complexes?

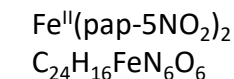
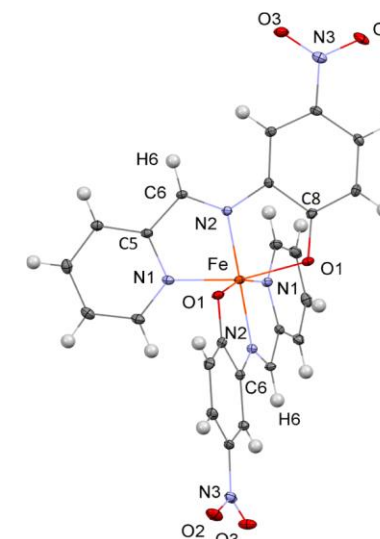
Material family	ΔS	ΔT	Reversibility	pressure	Thermal conductivity	Density	Cost and availability
Magnetostructural	**	*	**	*	****	****	**
Ferroelastic fluorides and oxyfluorides	**	**	***	***	**	**	****
Ferroelectrics	**	**	****	***	**	**	****
Superionic conductor	**	**	**	**	**	****	**
Hybrid organic–inorganic perovskites	**	***	****	****	*	*	**
Spin Crossovers	***	***	****	****	*	**	**
Polymers	***	***	***	*	**	*	****
Plastic crystals	****	****	**	**	**	*	****

*Poor, **Fair, ***Good, ****Excellent.

P. Lloveras and J. L. Tamarit, MRS Energy & Sustainability (2021) 8:3–15



J. Am. Chem. Soc. 2022, 144, 14, 6493–6503

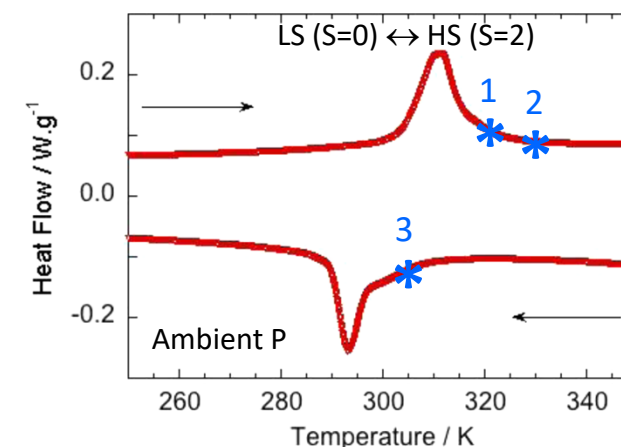


Good BC candidates:

- large $\Delta H = T\Delta S$ and large ΔV
- ΔT at low P
- T close to RT and small hysteresis

T_t heating = 312.4 K
 T_t cooling = 296.1 K
 $\Delta S = 30 \text{ J/mol} = 55.5 \text{ J/kg K}$
 Same C2/c monoclinic cell
 Volume change $\Delta V/V = 4.5\%$

O. lasco et al., Inorg. Chem. 2015, 54, 1791–1799



Temperature / K

Heat Flow / W.g⁻¹

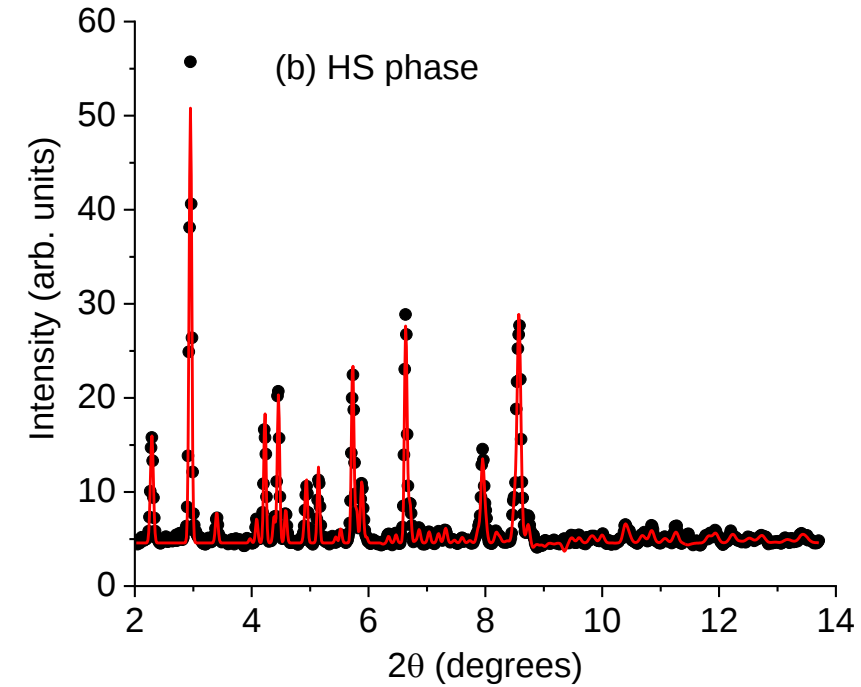
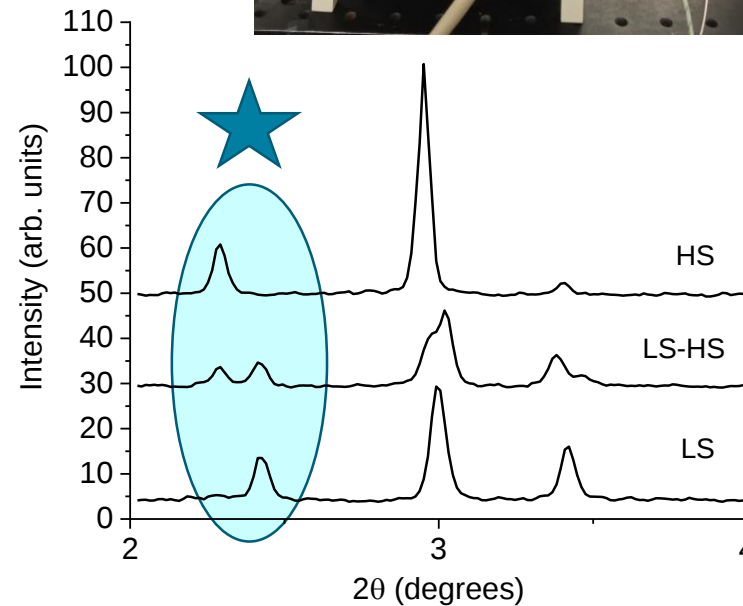
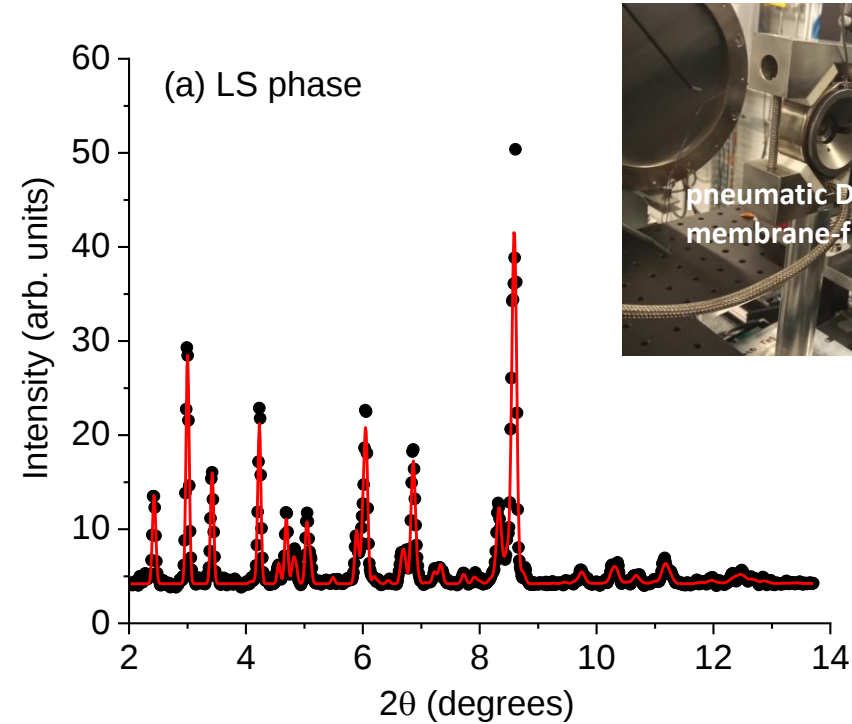
CSIC

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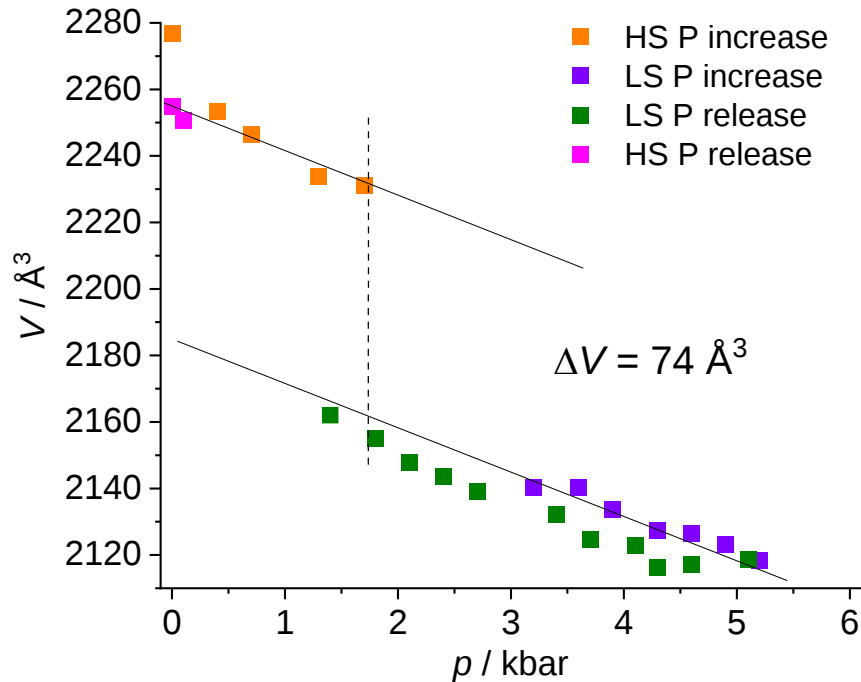
Results – Low pressure synchrotron powder X-ray diffraction

High-pressure endstation @ BL04-MSPD beamline at ALBA

E=29.22 keV; P range: AP to 6 kbar; isothermal HPXRD patterns at T=320 K → 330 K → 310 K

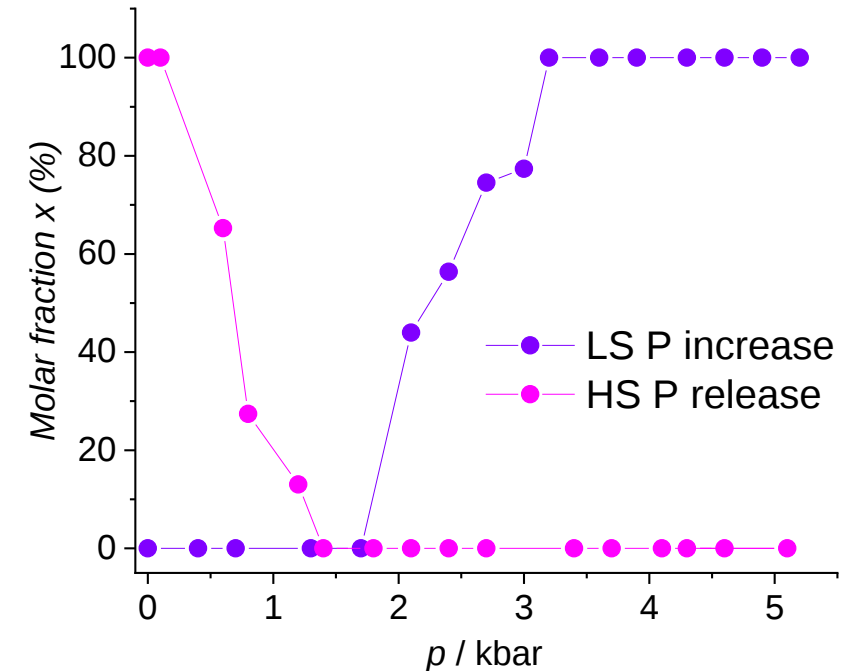


(1) T=320 K



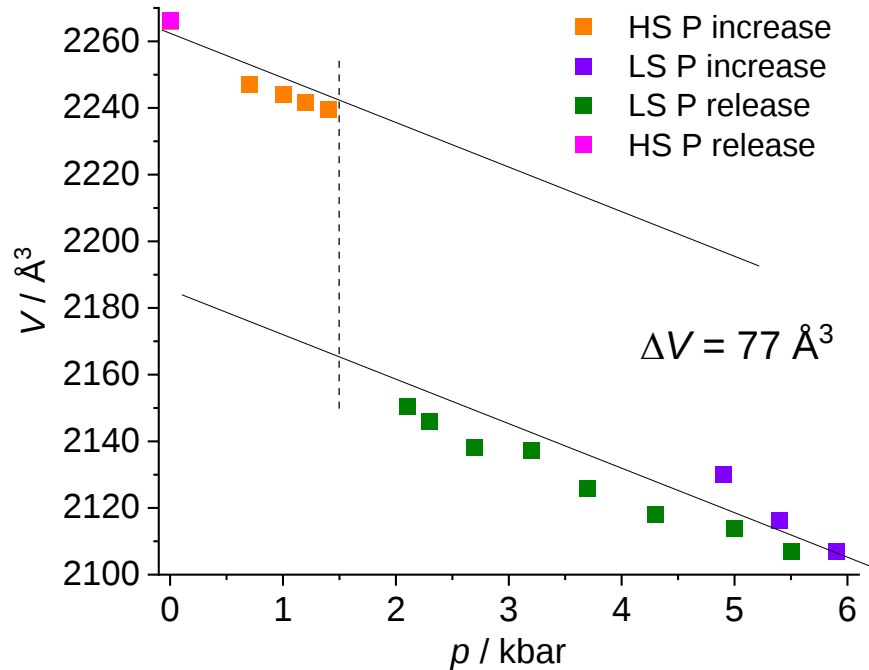
Isothermal compressibility of the LS phase:

$$\kappa = -\frac{1}{V} \left(\frac{\partial V}{\partial p} \right)_T = 0.05 \text{ GPa}^{-1}$$



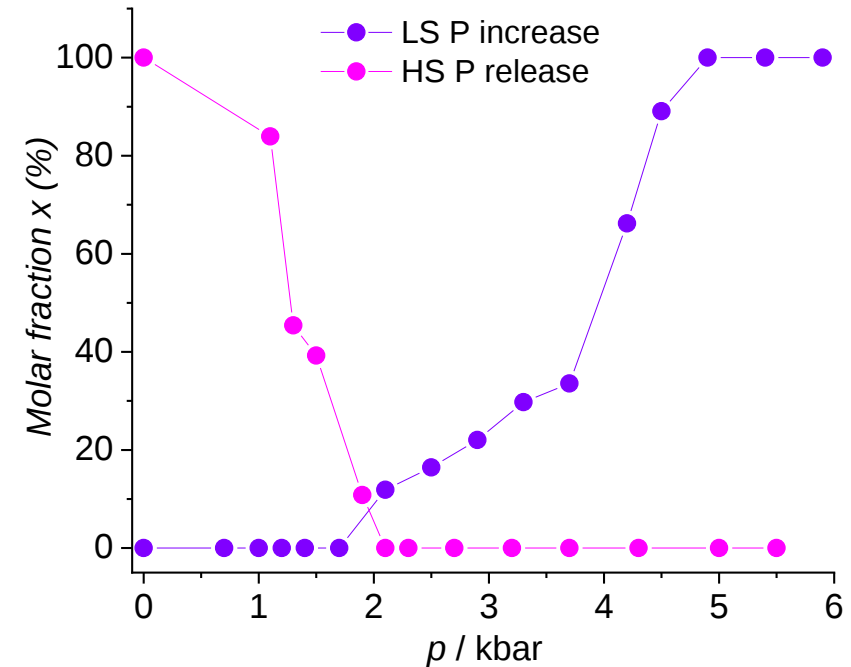
The LS phase starts at P=1.7 kbar on increasing P
& the HS phase is recovered at P=1.4 kbar on releasing P

(2) T=330 K



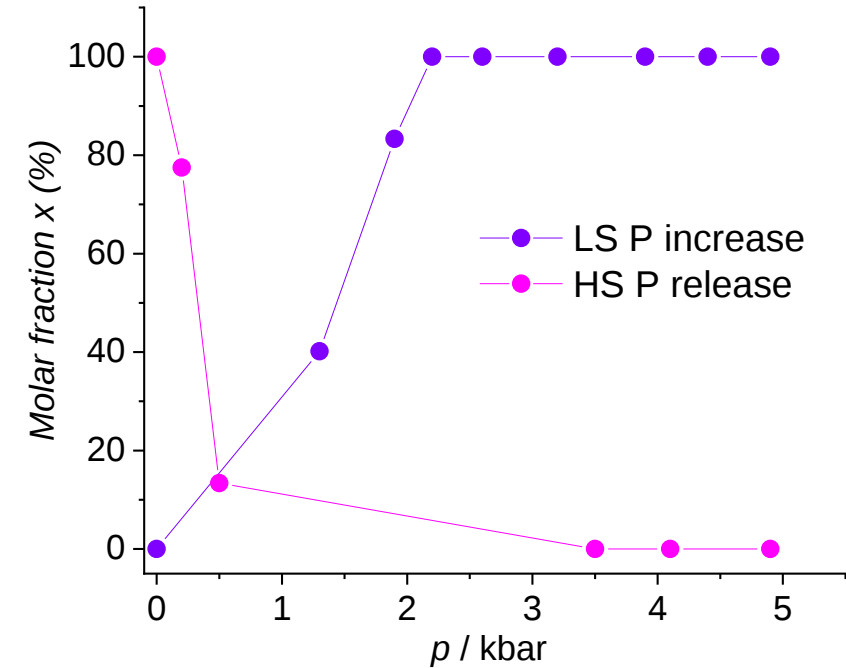
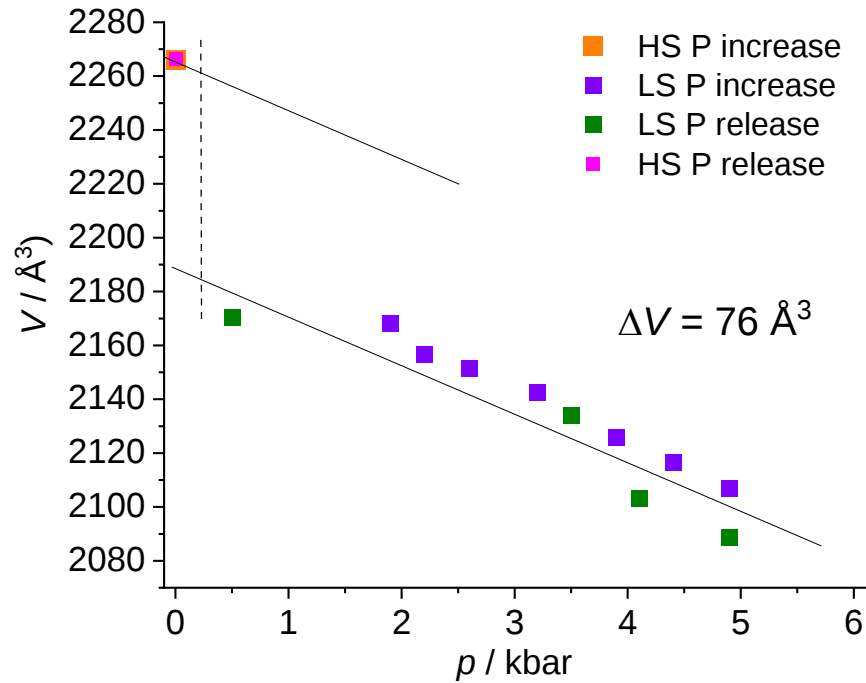
Isothermal compressibility of the LS phase:

$$\kappa = -\frac{1}{V} \left(\frac{\partial V}{\partial p} \right)_T = 0.048 \text{ GPa}^{-1}$$



The LS phase starts at P=1.7 kbar on increasing P
& the HS phase is recovered at P=1.9 kbar on releasing P

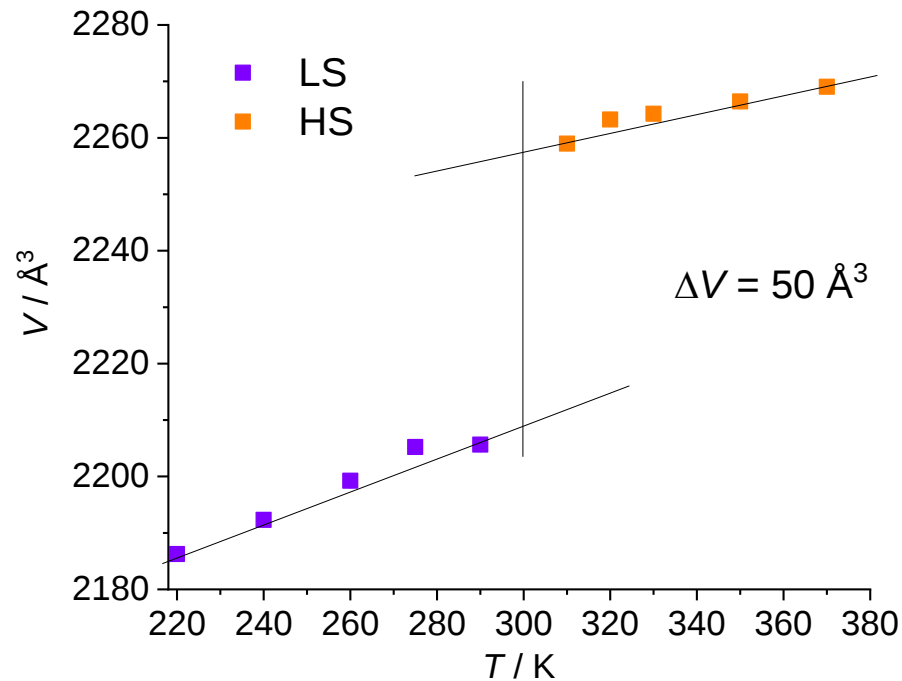
(3) T=310 K



Isothermal compressibility of the LS phase:

$$\kappa = -\frac{1}{V} \left(\frac{\partial V}{\partial p} \right)_T = 0.085 \text{ GPa}^{-1}$$

The LS phase starts at ~1 kbar on increasing P
& the HS phase is recovered at 0.5 kbar on releasing P



Volumetric expansion (LS): $\beta = \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_p = 1.63 \times 10^{-4} \text{ K}^{-1}$

Volumetric expansion (HS): $\beta = \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_p = 6.78 \times 10^{-5} \text{ K}^{-1}$

Estimation of the lattice entropy change at the transition:

$$\Delta S_v = (\beta / \kappa) \Delta V$$

κ average = $0.06 \text{ GPa}^{-1} = 6 \times 10^{-11} \text{ Pa}^{-1}$.

β average = $1.1 \times 10^{-4} \text{ K}^{-1}$.

ΔV isothermal = $50 \text{ Å}^3 = 1.4 \times 10^{-5} \text{ m}^3 \text{ kg}^{-1}$

Lattice entropy change: $\Delta S_v = 25.7 \text{ J kg}^{-1} \text{ K}^{-1}$

Magnetic entropy change: $\Delta S_m = R \ln(5) = 24.8 \text{ J kg}^{-1} \text{ K}^{-1}$

$$\Delta S_t = \Delta S_v + \Delta S_m = 50.5 \text{ J kg}^{-1} \text{ K}^{-1}$$



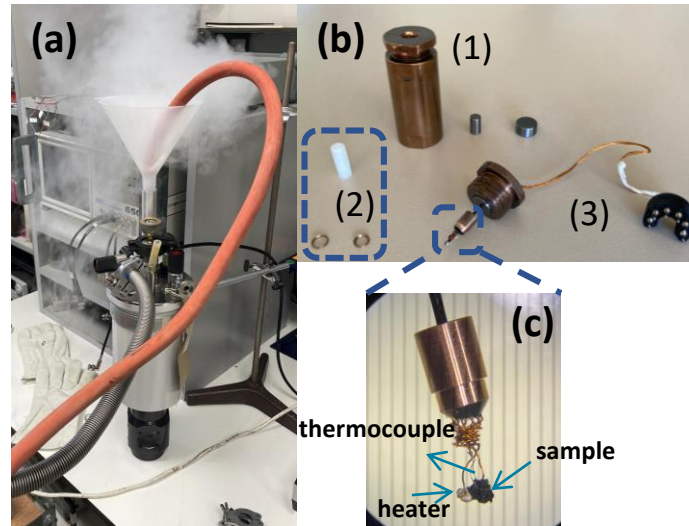
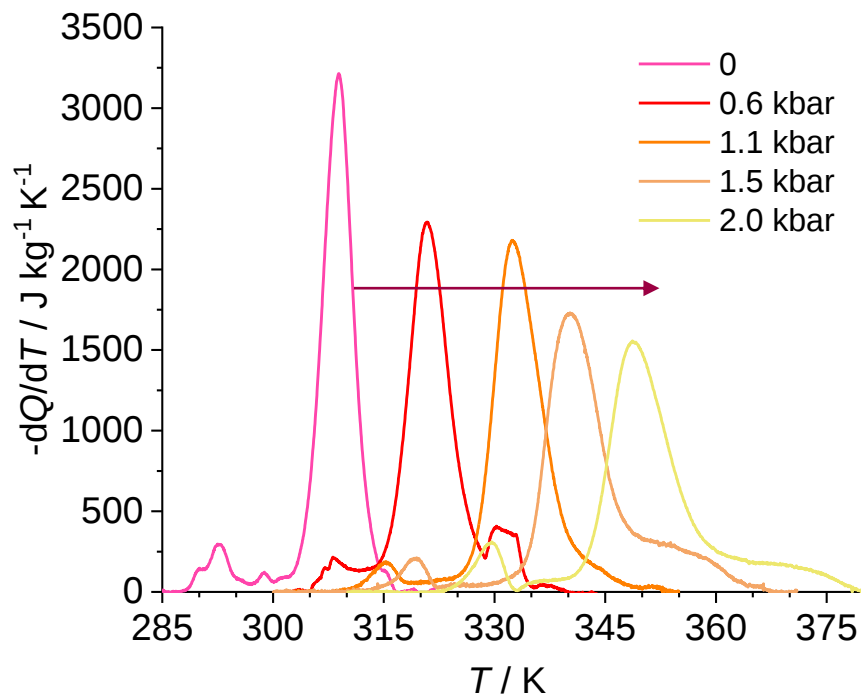
$$\Delta S_t = 55.5 \text{ J kg}^{-1} \text{ K}^{-1}$$



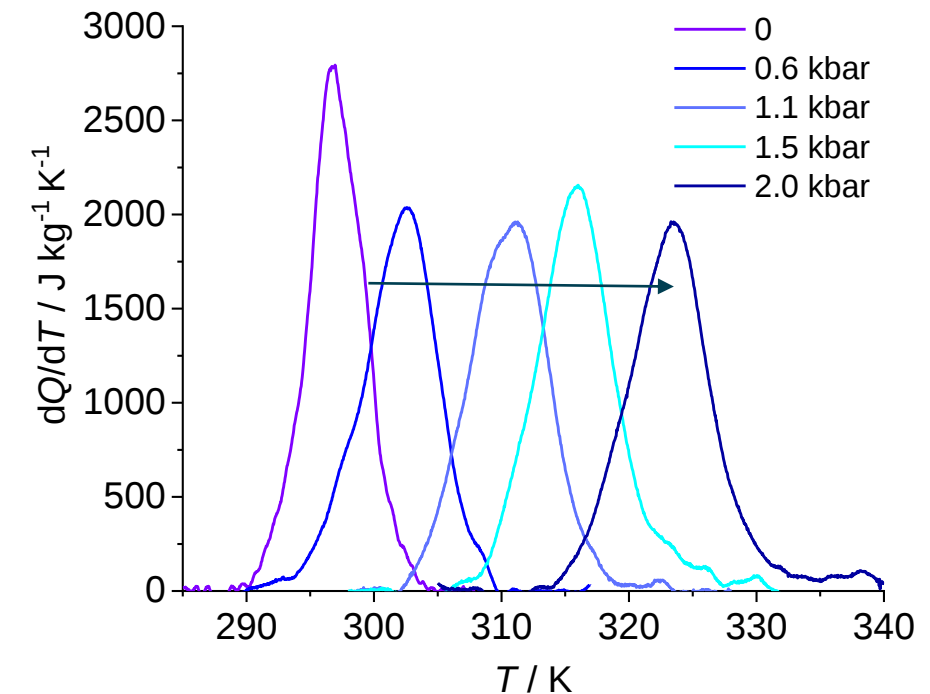
O. Iasco et al., Inorg. Chem. 2015, 54, 1791–1799

Heat capacity anomaly as a function of pressure

Heating



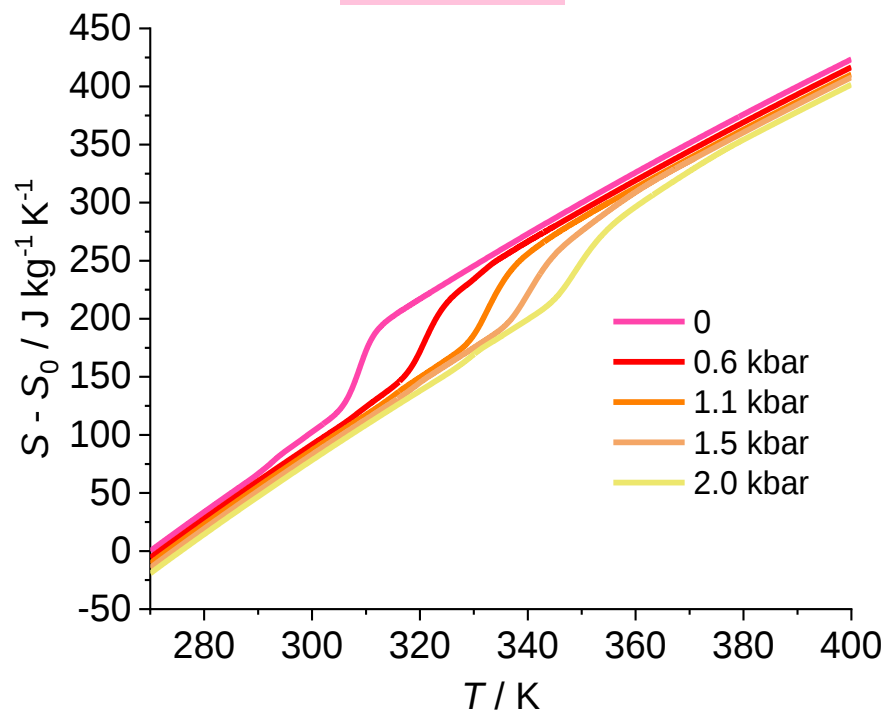
Cooling



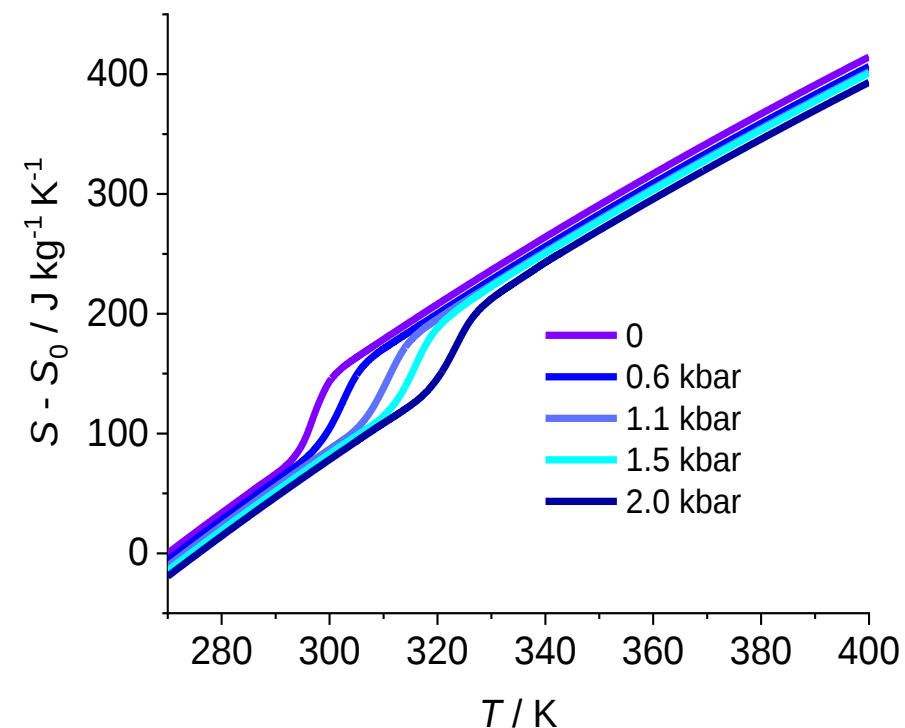
Entropy curves as a function of temperature at different P values

$$S(T, p) - S_0(T_0, p_0) = - \left(\frac{\partial V}{\partial T} \right)_{p=p_0} \Delta p + \int_{T_0}^T \frac{1}{T} \left(c_p + \frac{dQ}{dT} \right) dT \quad T_0 = 270 \text{ K}$$

Heating



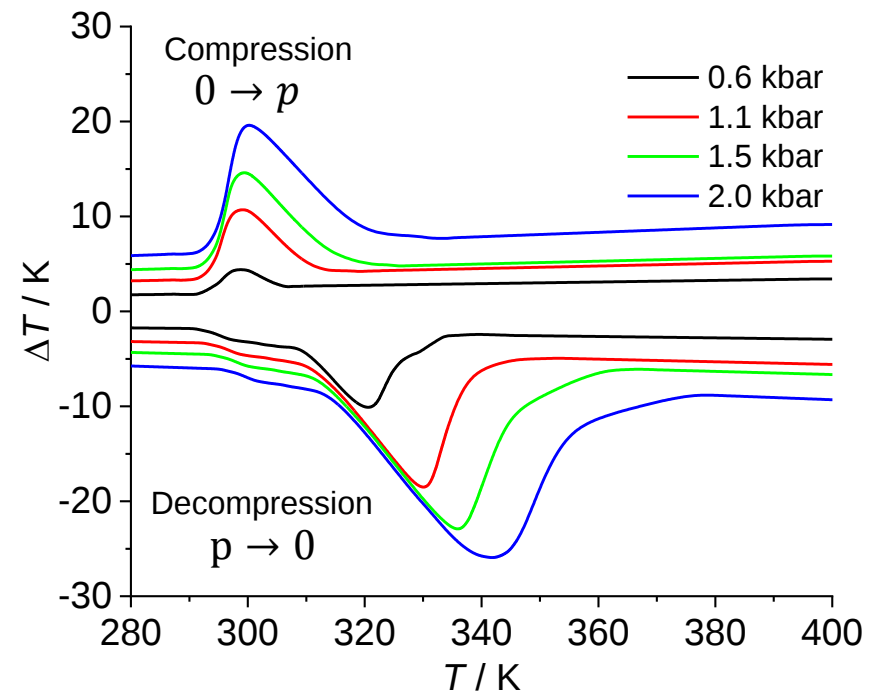
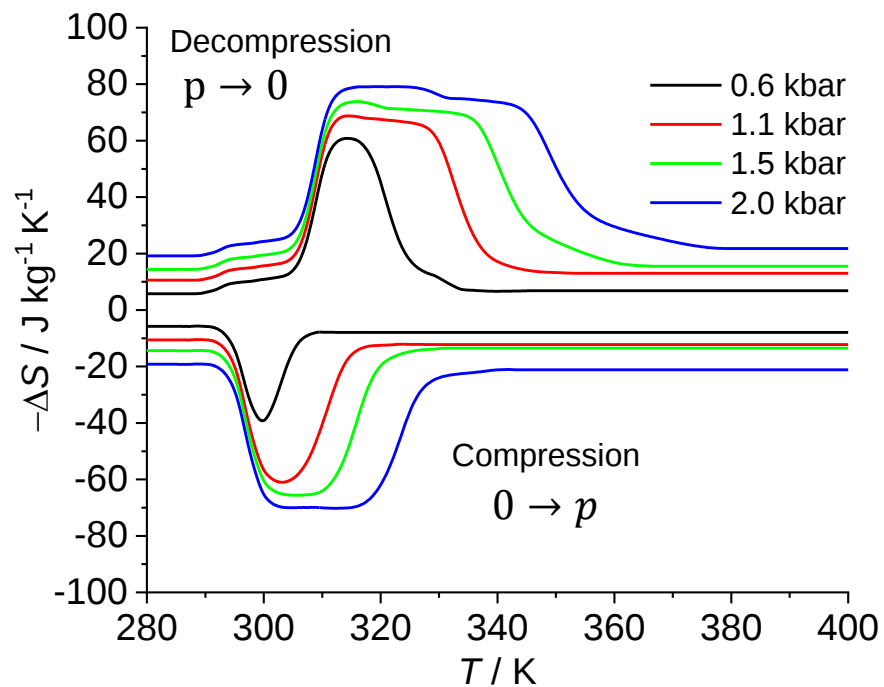
Cooling



Isothermal Entropy change

Adiabatic Temperature change

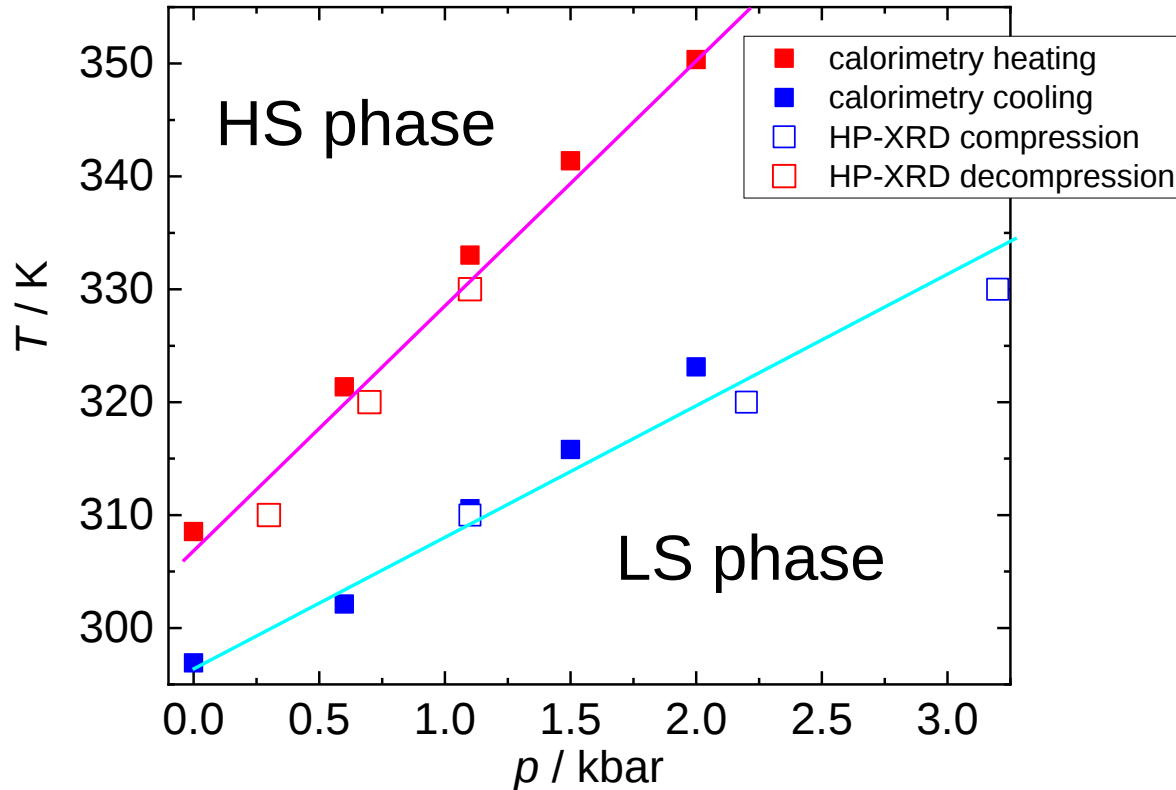
induced by compression (cooling runs) or decompression (heating runs)



For low values of p (0.6 kbar), significantly large $\Delta S \approx 50 \text{ J Kg}^{-1} \text{K}^{-1}$ and $\Delta T = 7 \text{ K}$ are found and they increase up to $\Delta S = 70 \text{ J Kg}^{-1} \text{K}^{-1}$ and $\Delta T = 20 \text{ K}$ for $p = 2 \text{ kbar}$.



$T - P$ phase diagram



□ $\text{Fe}^{\text{II}}(\text{pap-5NO}_2)_2$ is a good candidate for BC effect as T_t spans from 297 K at ambient pressure to 330 K at 3 kbar.

□ A remarkable large BC effect with $\Delta S = 70 \text{ J Kg}^{-1} \text{ K}^{-1}$ and $\Delta T = 20 \text{ K}$ for $p = 2 \text{ kbar}$ is reported.



□ Estimated transition temperature displacement with pressure:

(1) LS to HS transition

(red symbols/magenta line): $dT_t/dp = 22.4 \text{ K kbar}^{-1}$

agrees with Clausius-Clapeyron equation

(2) HS to LS transition

(blue symbols/cyan line): $dT_t/dp = 11.1 \text{ K kbar}^{-1}$

The hysteresis increases with pressure





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David Gracia
Vera Cuartero
Emmanouil Charkiolakis
Marco Evangelisti
Javier Blasco

Université Paris-Saclay

Marie-Laure Boillot
Talal Mallah

BL04 – MSPD beamline

Catalin Popescu

Thank you for your attention!